Rare Decays: The Sharpest Tools for New Physics at LHCb

Debashis Sahoo



October 2, 2025

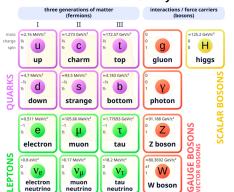


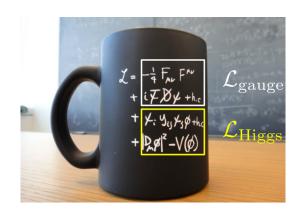




Current Understanding

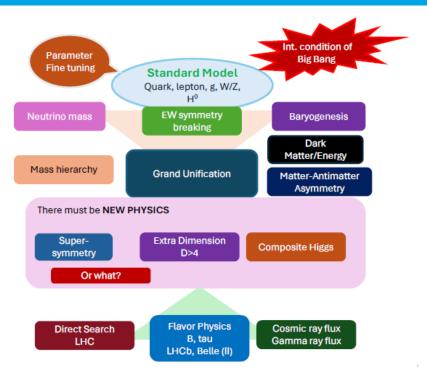
Standard Model of Elementary Particles





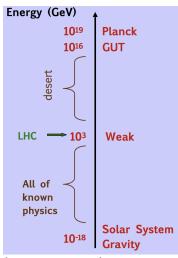
- Fundamental Interactions:
 - Electromagnetic \rightarrow photon (γ) as carrier
 - Strong \rightarrow gluon (g) as the carrier
 - Weak $\to W^{\pm}, Z^0$ as the carrier
 - Gravity \rightarrow Graviton (?) as the carrier (Not included in the SM)
- $\mathcal{L}_{\text{gauge}} \rightarrow$ mathematically elegant symmetries $\mathcal{L}_{\text{Higgs}} \rightarrow$ ad hoc Yukawas

Beyond the Standard Model



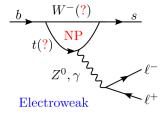
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Flavor as a discovery tool



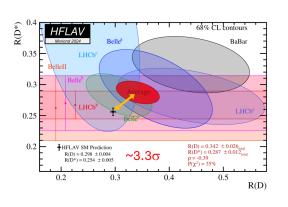
[J. Hewett, LISHEP09]

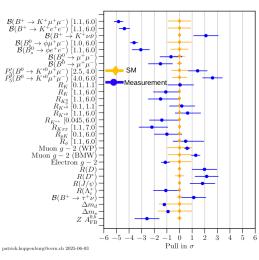
- Long history of flavor as an "indirect" probe for new heavy particles:
 - weak nuclear β decay \Rightarrow heavy W/Z
 - $-K_L^0 \to \text{GIM suppression} \Rightarrow \text{charm}$
 - $-B^0$ -mixing at ARGUS \Rightarrow heavy top
 - SM-like $\mathcal{B}(B_s^0 \to)$ at the LHC \Rightarrow tight limits on MSSM/SUSY



• Even if $\Lambda_{\rm NP} \gg {
m TeV}$, precision flavor can probe the "desert" via rare loop-mediated processes.

Recent Flavor Anomalies

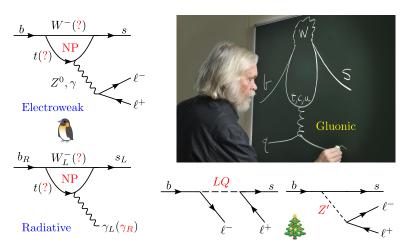




- Tension is at 3.3σ in $b \to c\ell\nu$ transitions
- BR measurements differ from predictions

Rare b-decays

- $b \to s(d)$ flavor changing neutral currents: loop-suppressed in SM.
- New Physics (NP) can enter both at loop- and tree-levels.
- These kinds of rare decays have a long tradition of discovering new things even before doing the actual observations.



EFT tools for rare decays

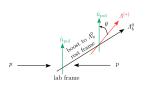
- Renormalizability requires the \mathcal{L}_{SM} to have dim $d \leq 4$ operators.
- Eg.: $m_{\phi}^2 \phi^2$, $m_{\psi} \overline{\psi} \psi \Rightarrow (m_{\phi}/E)^2$, (m_{ψ}/E) UV-safe behavior.
- We can include d > 4 operators if we regard the SM as an low energy effective theory. Comes with a cutoff scale, Λ .

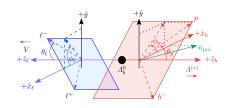
$$\mathcal{L}_{\text{eff}}(x) = \mathcal{L}_{\text{SM}}(x) + \sum_{d>4} \frac{C_i}{\Lambda^{d-4}} \mathcal{O}_i^{(d)}(x)$$
local operators

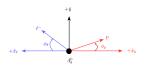
- Amplitudes will have $(E/\Lambda)^{d-4}$ behavior: bad at high-E, but suppressed at $E \ll \Lambda$. Access to heavy $(\Lambda_{\rm NP})$ fields from NP.
- Relevant for RD: d=6 operators. $\mathcal{A}_{\text{eff}} \sim \frac{C_{\text{SM}}}{m_W^2} + \frac{C_{\text{NP}}}{\Lambda_{\text{NP}}^2}$.

Angular analyses as a tool for NP searches

- Huge LHC statistics allow precision measurements of angular observables in $b \to s \ell^+ \ell^-$ and $b \to s \vec{\gamma}$. Direct access to $C_i^{\rm NP}$.
- Eg., $\vec{\Lambda_b^0} \equiv |[ud]\vec{b}\rangle$ reflects the properties of the *b*-quark, with [ud] as spectator diquark.
- $\bullet \ \vec{\Lambda_b^0} \to \Lambda^*(\to pK^-)\ell^+\ell^- \colon \{ {\bf q^2} \equiv m_{\ell^+\ell^-}^2, {\bf k^2} \equiv m_{pK}^2 \} \, + \, \{ {\theta_\ell, \theta_p, \phi_\ell, \phi_p} \}$



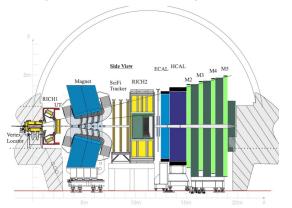


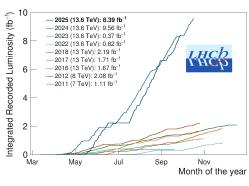


• If Λ_b^0 is unpolarized, $\phi_\ell = 0$, $\chi \equiv \phi_p$. Similar definitions for $B^0 \to K^+\pi^-\ell^+\ell^-$, $B_s^0 \to K^+K^-\ell^+\ell^-$.

LHCb: a discovery tool

[2024 JINST 19 P05065]





- ullet Single forward-arm spectrometer dedicated to c- and b-physics.
- Major Upgrade-I installed during LS2. Copious Run 3 data taking-ongoing.
- In this talk, I will focus on two rare decay analyses:
 - Electroweak: $\Lambda_b^0 \to pK^-\ell^+\ell^-$ using Run 1+2 (2011-18) data
 - Radiative: triggers and $B_c^+ \to D_s^{*+} (\to D_s^+ \gamma) \gamma$ in Run 3 (2024)

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Lepton flavor Universality (LFU) in SM

 \bullet LFU in the SM: Electroweak couplings of gauge bosons to leptons are independent of their flavor \to

$$\Gamma(z \sim \sqrt{\frac{e}{e}}) = \Gamma(z \sim \sqrt{\frac{\mu}{\mu}})$$

$$\Gamma(w^{\dagger} \sim \sqrt{\frac{e}{\bar{\nu}_e}}) = \Gamma(w^{\dagger} \sim \sqrt{\frac{\mu}{\bar{\nu}_{\mu}}})$$

- LHCb is the LFU test industry,
- Thrust has been the ratios:

$$R(H_c) = \frac{\mathcal{B}(H_b \to H_c \tau \nu_{\tau})}{\mathcal{B}(H_b \to H_c \mu \nu_{\mu})}, \ R(H_s) = \frac{\mathcal{B}(H_b \to H_s \mu \mu)}{\mathcal{B}(H_b \to H_s e e)}$$

$$\text{ere } H_b \in \{B^0 \ B^+, \ \Lambda^0 \ B^0\} \ H_s \in \{K^{(*)} \ \phi \ \Lambda^{(*)}\}$$

Where
$$H_b \in \{B^0, B_{(c)}^+, \Lambda_b^0, B_s^0\}, H_s \in \{K^{(*)}, \phi, \Lambda^{(*)}\},$$

 $H_c \in \{D^{(*)+,0}, D_s^{(*)+}, \Lambda_c^{(*)-}, J/\psi\}$

LFUV via angular analysis of $\Lambda_{\rm b}^0 o { m pK}^-\ell^+\ell^-$ (Ongoing)

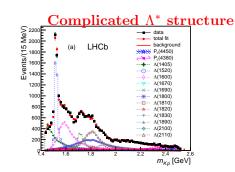
Phys. Rev. Lett. 115, 072001 (2015)

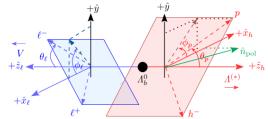
- Wide overlapping resonances is a major complication
- The angular rate for $\Lambda_b^0 \to p K^- \ell^+ \ell^-$

$$\frac{d^4\Gamma}{d\mathbf{q^2}\,d\cos\theta_\ell\,d\cos\theta_\mathbf{p}\,d\chi} = \sum_{\mathbf{i}} \mathbf{I_i}(\mathbf{q^2}) \mathbf{f_i}(\theta_\ell,\theta_\mathbf{p},\phi)$$

• Measure the LFUV observables

$$S_i \equiv \tilde{I}_{i,\mu\mu} - \tilde{I}_{i,ee}$$



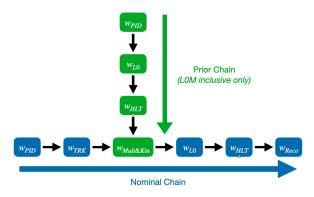


Run 1+2 Analysis Chain $\Lambda_{\rm b}^0 \to { m pK}^-\ell^+\ell^-$ (Ongoing)

- Run 1+2 Data and MC samples \checkmark
- Online-Selections ✓
- \bullet Trigger and Pre-BDT selections \checkmark
- Simulation correction Chain $\sqrt{}$ \leftarrow will talk from here
- \bullet Combinatorial background suppression by BDT \checkmark
- Fit to Invariant Λ_b^0 mass (Ongoing)
- Perform Angular analysis
- Systematics calculations and results

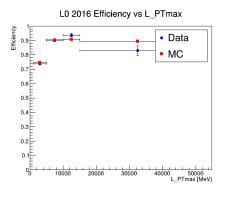
Corrections to the simulation

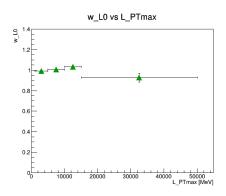
- Data-simulation differences in trigger, kinematic and PID
- Corrections done in two "chains": trigger corrections before (prior) or after (nominal) kinematic corrections.



• Following the strategy of other LFUV LHCb analyses, R_{K,K^*} and $R_{K\pi\pi}$.

- Example: lepton p_T dependence of the L0 trigger corrections
- Prior chain: $w_{L0} = \frac{\epsilon_{data}}{\epsilon_{MC}}$

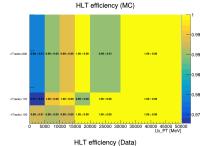


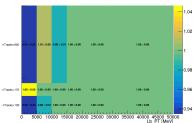


Corrections to the simulation: HLT

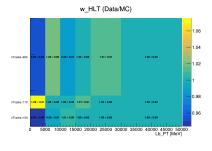
LHCb Unofficial

• Prior chain:
$$w_{HLT} = \frac{\epsilon_{data}}{\epsilon_{MC}}$$





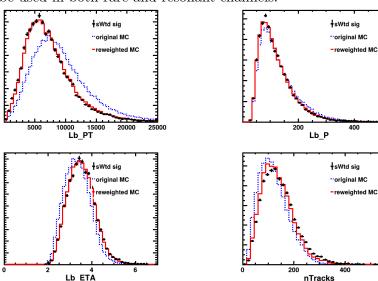
• 2D dependence of weights in bins of Λ_b^0 p_T and track multiplicty.



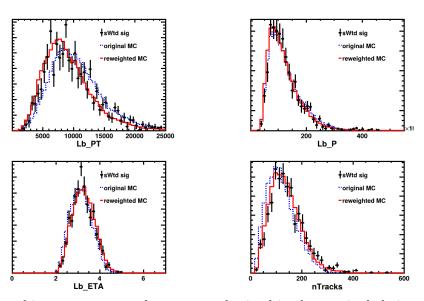
Kinematic variable corrections, $w_{Multi\&kin}$

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- Use the previously calculated weights
- Correction by a Gradient Boost reweighter
- Separately for each trigger category. Corrections from the resonant channels will be used in both rare and resonant channels.



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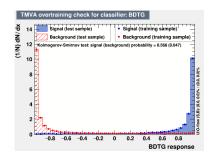


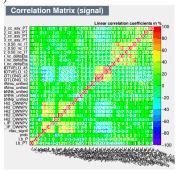
• After this, w_{L0} , w_{HLT} and w_{Reco} are obtained in the nominal chain.

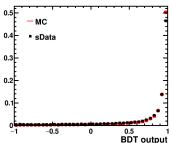
BDT training for $\Lambda_b^0 \to pK^-J/\psi(\to \mu^+\mu^-)$

• BDT trained on corrected simulation (signal) and upper-mass sideband $m(pK^-\mu^+\mu^-) > 5800$ data (background proxy).

- BDT includes a large number of kinematic, vertexing, PID, and isolation variables
- Excellent performance and data-simulation agreement







BDT training for $\Lambda_b^0 \to pK^-J/\psi(\to e^+e^-)$

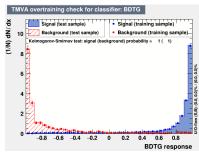
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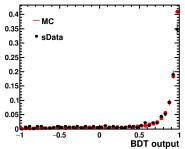
• For ee, due to low statistics, BDT with k-folding and cross-validation

• Good separation and no overtraining despite a large number of variables in the BDT

 Good agreement between corrected simulation and background-subtracted data

• All correlations are preserved between data and simulations

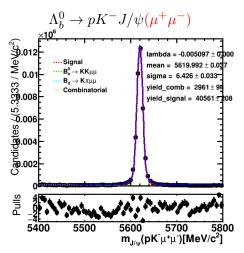




Preliminary resonant channel mass peaks

LHCb Unofficial

• Checks for Run 2 (2016) for one particular trigger (TOS)



 $\Lambda_b^0 \to pK^-J/\psi(e^+e^-)$ Candidates / (9.4118 / MeV<u>/</u>c² 00 00 00 00 00 mean = 5620.27 + 0.27 slope = -0.003469 ± 0.0004 yield_comb = 1002 ± 122 vield lb2pkeepi0 = 65 ± 9 Combinatoria vield signal = 3999 ± 75 5400 5600 5800 6000 $m_{1/w}(pK^{-}e^{+}e^{-})[GeV/c^{2}]$

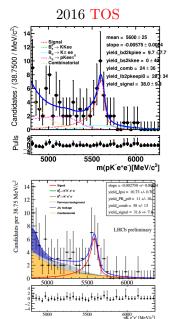
Signal yield increased by a factor of 2.3 compared to published R_{nK}

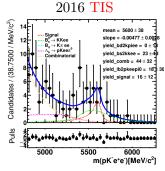
Signal yield increased by a factor of 1.5 compared to published R_{nK}

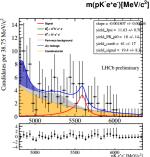
• Need to understand the large improvement.

Preliminary $\Lambda_b^0 \to pK^-e^+e^-$ fits for 2016 data

LHCb Unofficial







- Showing Run 2 2016 comparisons for different triggers
- Good agreement with the published 2020 R_{pK} paper
- Signal yields are of similar order with the same background
- Ongoing work: better constrain the background components

Ongoing...

- Analysis note is being written at the same time
- Finalizing the mass fits for all the years and rare decays
- Next, perform the angular analysis
- We are working on a phenomenology paper to get the angular observables for any Λ spin.

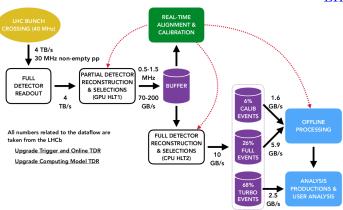
Run 3

LHCb Run 3 Data Flow

[LHCB-TDR-016]

• Level0 hardware trigger bottleneck removed.

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- First time completely software trigger in a high energy experiment!
- I contributed significantly to the HLT2 trigger, which performs the more computationally expensive full reconstruction online during data taking for a specific physics analysis.

Radiative Exclusive HLT2 trigger

LHCb/gitlab/Moore/link

- Built strong radiative program at LHCb in Run 3:
- Designed the following Physics program through HLT2 trigger:

 LHCb Unofficial

$$\bullet B^0_s \to K^+K^-\gamma$$

$$\bullet$$
 $B^0 \to K^+\pi^-\gamma$

$$\bullet \ B^+ \to K^+ \pi^+ \pi^- \gamma$$

$$\bullet$$
 $B^+ \to K^+ \phi \gamma$

$$\bullet \ B^0 \to K_s^0 \pi^+ \pi^- \gamma$$

$$\bullet \ B^0 \to K^0_s \phi \gamma$$

$$b^+ \to K^+ \omega (\pi^+ \pi^- \pi^0) \gamma$$

$$B^0 \to K_s^0 \omega(\pi^+\pi^-\pi^0) \gamma$$

$$\bullet \ B_s^0 \to K_s^0 K^+ \pi^- \gamma$$

$$\bullet \ \Lambda_b \to K_s^0 p \pi^- \gamma$$

•
$$\Lambda_b \to pK^-\gamma$$

$$\bullet \ B^0 \to \pi^+\pi^-\gamma$$

•
$$B_c^+ \to D^*(2010)^+ \gamma$$

$$\bullet \ B_c^+ \to D_s^{*+} \gamma$$

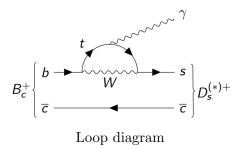
$$\begin{array}{c} \bullet \ B^+ \rightarrow \\ K^+ \eta (\pi^+ \pi^- \pi^0) \gamma \end{array}$$

•
$$B^0 \to K_s^0 \eta(\pi^+\pi^-\pi^0) \gamma$$

• The data collected with these triggers will lead to at least 10 publications.

Radiative B_c^+ decays

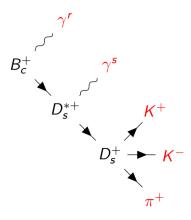
• B_c meson contains two heavy quarks, b and $c \to \text{Large number of decay modes}$.



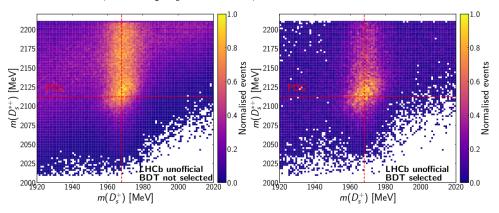
- B_c can decay radiatively as $B_c^+ \to D_s^+ \gamma$ and $B_c^+ \to D_s^{*+}[D_s^+ \gamma] \gamma$
- Together with the EPFL group, we have started looking at $B_c^+ \to D_s^{*+} \gamma$ decays.
- Never Searched experimentally
- Theoretically expected BF:~ $\mathcal{O}(10^{-5})$

Rare $B_c^+ \to D_s^{*+}[D_s^+\gamma]\gamma$

- Only possible at LHCb
- Opportunity to study radiative transitions in the B_c sector
- Many B_c^+ in Run 3
- Challenges:
 - D_s flies, B_c vertex is not reconstructible. Similar to $\Lambda_b^0 \to \Lambda \gamma$
 - 2 photons final state
 - 1 soft photon (γ^s); introduces lots of background, photon identification challenging at low PT



• Samuël Bakker, EPFL project student, looked at 2024 data



- There are D_s^{*+} and D_s^{+} mass peaks in data.
- Next, we will study the misidentified and peaking backgrounds

Neutral PID calibration at LHCb

LHCb Unofficial LHCb-INT-2025-002

- Identification of photons and neutral pions at LHCb is key to several physics analyses.
- Designed LHCb Calo. triggers.
- Decay channels used (large branching fraction) for calibration:

$$B_{(s)}^{0} \to K^{*0}/\phi \gamma$$

$$D^{*+} \to D^{0}[K^{+}\pi^{-} \pi^{0}]\pi^{+}$$

$$D_{s}^{*+} \to D_{s}^{+} \gamma$$

$$\eta \to \mu^{+}\mu^{-} \gamma$$

$$D_{(s)}^{+} \to \eta'[\pi^{+}\pi^{-} \gamma]\pi^{+}$$

• Calibration will be used by the whole LHCb collaboration

Future Plans

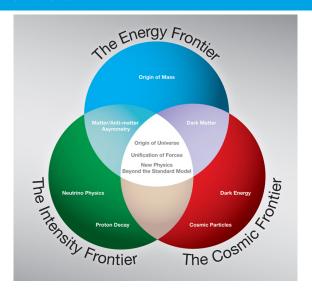
- The journey at LHCb has been exciting, received the YSF grant at La Thuile, Italy, for EMTF work
- Recently won the EKOP scholarship in Hungary (additional support to exceptional researchers)
- After finishing the $\Lambda_b^0 \to p K^- \ell^+ \ell^-$ analysis, I plan to focus on radiative decays
- First searches of decays like $B_c^+ \to D_s^{*+} \gamma$, $\Lambda_b \to K_s^0 p \pi^- \gamma$, $B^0 \to K_s^0 \pi^+ \pi^- \gamma$ would be the immediate thrust
- Observation of these decays will open new avenues for TDCPV, photon polarization study
- I would also be interested in contributing to LHCb Upgrade II.

Conclusion

- Study of rare decays is absolutely crucial for new Physics search
- Designed RD HLT2 trigger lines and also added new Semileptonic and B2OC trigger lines
- Also, contributed as RD RTA liaison for 2025 data taking
- $\Lambda_b^0 \to p K^- \ell^+ \ell^-$, planning to enter into RC by the end of this year

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The Three Frontiers



- At the intensity frontier, rare decays of subatomic particles like b meson, d meson, τ lepton are studied. Their probability of occurrence is very small.
- We compensate for that by having high luminosity, a large number of these particles being produced.

Reduce Combinatorial backgrounds by BDT

List of variables used for BDT training:

- $p_T(\Lambda_b^0)$
- $p(\Lambda_b^0)$
- χ^2_{DTF} prob
- ctau_signif
- k_IP_OWNPV, p_IP_OWNPV
- Lb_IPCHI2_OWNPV, p_IPCHI2_OWNPV, k_IPCHI2_OWNPV, L1_IPCHI2_OWNPV,L2_IPCHI2_OWNPV
- p_ProbNNp, p_ProbNNk, k_ProbNNk
- L1_ProbNNmu, L2_ProbNNmu,

Isolation variables:

- Lb_TRKISOBDTLONG_12, Lb_TRKISOBDTLONG_45
- Lb_TRKISOBDTVELO_12, Lb_TRKISOBDTVELO_45
- Lb_L2_0.50_nc_deltaEta, Lb_Kaon_0.50_nc_deltaEta, Lb_Proton_0.50_nc_deltaEta
- Lb_L2_0.50_nc_IT, Lb_Proton_0.50_nc_IT, Lb_Kaon_0.50_nc_IT

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Triggers

Mutually exclusive L0-Hardware Triggers:

- TIS: LOHadron_TIS (Lb) || LOMuon_TIS (Lb) || LOElectron_TIS (Lb)
- MTOS: LOMuon_TOS $(\mu 1, \mu 2)$ && !TIS(Lb)
- ullet ETOS: LOElectron_TOS (e1,e2) && !TIS(Lb)

After that, further software-level triggers HLT1 and HLT2 are being applied.

Basis of local operators for $b \rightarrow s$ penguins

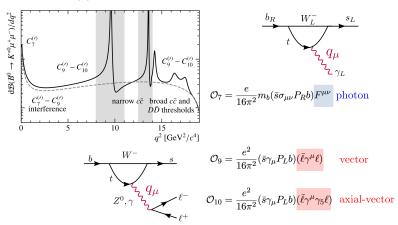
• (V - A) LH operators consistent with SM symmetries:

$$\begin{split} \mathcal{O}_1^u &= (\bar{s}\gamma_\mu T^a P_L u) \, (\bar{u}\gamma^\mu T^a P_L b) \\ \mathcal{O}_2^u &= (\bar{s}\gamma_\mu P_L u) \, (\bar{u}\gamma^\mu P_L b) \\ \mathcal{O}_2^u &= (\bar{s}\gamma_\mu P_L u) \, (\bar{u}\gamma^\mu P_L b) \\ \mathcal{O}_2^u &= (\bar{s}\gamma_\mu P_L u) \, (\bar{u}\gamma^\mu P_L b) \\ \mathcal{O}_1^c &= (\bar{s}\gamma_\mu T^a P_L c) \, (\bar{c}\gamma^\mu T^a P_L b) \\ \mathcal{O}_2^c &= (\bar{s}\gamma_\mu P_L c) \, (\bar{c}\gamma^\mu P_L b) \\ \mathcal{O}_2^c &= (\bar{s}\gamma_\mu P_L c) \, (\bar{c}\gamma^\mu P_L b) \\ \mathcal{O}_3^c &= (\bar{s}\gamma_\mu P_L b) \sum_q (\bar{q}\gamma^\mu q) \\ \mathcal{O}_4^c &= (\bar{s}\gamma_\mu T^a P_L b) \sum_q (\bar{q}\gamma^\mu T^a q) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{s}\gamma_\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &= \frac{e^2}{16\pi^2} (\bar{r}\gamma^\mu P_L b) (\bar{\ell}\gamma^\mu \gamma_5 \ell) \\ \mathcal{O}_{10}^c &=$$

• $\mathcal{O}_{1,2}$ (4-quark tree), \mathcal{O}_{3-6} (4-quark penguins), \mathcal{O}_8 (gluon penguin)

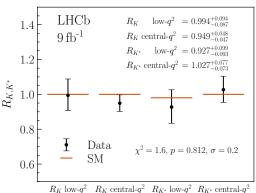
The three dominant contributions

• The dominant $\mathcal{O}_{7,9,10}$ contributions, as a function of q^2 :



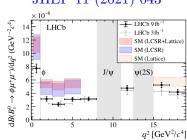
• The primed terms are the RH (quark) operators, suppressed in the SM, but can be enhanced in NP scenarios.

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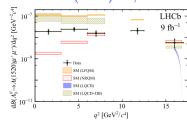


Compatible to the SM

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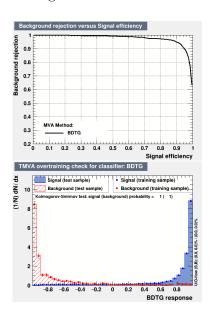
PRL 131 (2023) 15, 151801

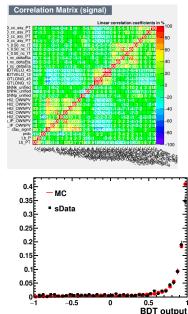


BR measurements differ from predictions

BDT training for $\Lambda_b^0 \to pK^-J/\psi(e^+e^-)$

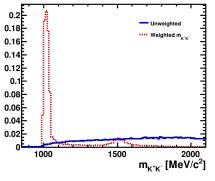
• For the electron case, due to low statistics, BDT k-folding with cross-validation is being used.

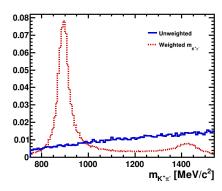




October 2, 2025

Background Shapes in MC





- Background samples were generated with flat in m_{KK} and $m_{K\pi}$.
- m_{KK} and $m_{K\pi}$ shapes are corrected.
- Get the expected yield =