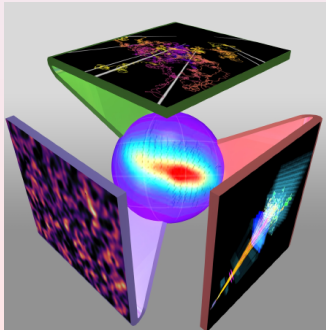


SFB 1491

Highlights from the Astro-Theory Projects

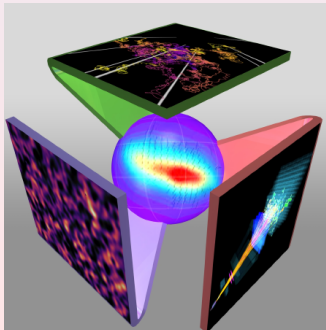
Horst Fichtner



SFB 1491

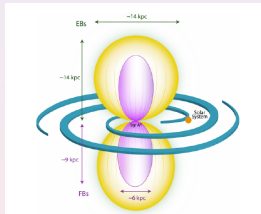
Highlighting the Astro-Theory Projects

Horst Fichtner



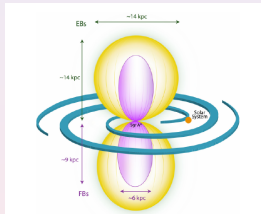
The Astro Projects

A1: Multimessenger signatures of Gal. CR transport

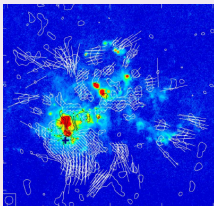


The Astro Projects

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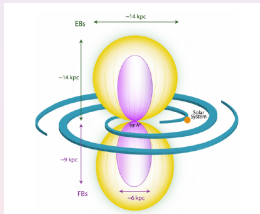


A2: CR signatures in dwarf galaxies

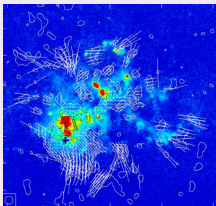


The Astro Projects

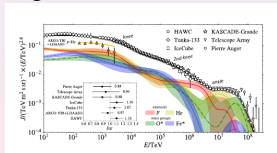
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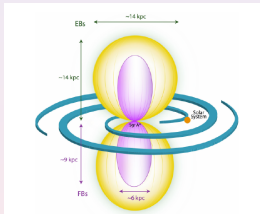


A3: CR transport in the transition region from Galactic to extragalactic origin

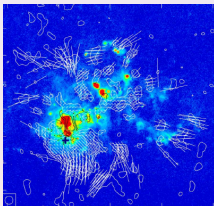


The Astro Projects

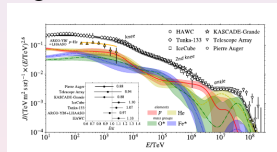
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A3: CR transport in the transition region from Galactic to extragalactic origin

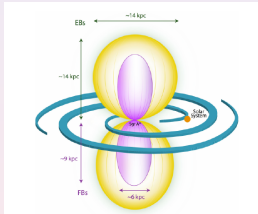


A4: Magnetohydrodynamical halos of starforming galaxies

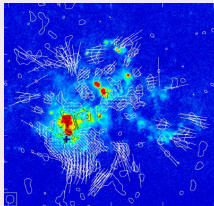


The Astro Projects

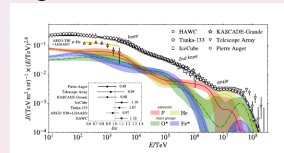
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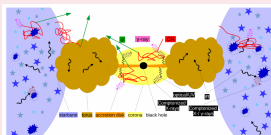
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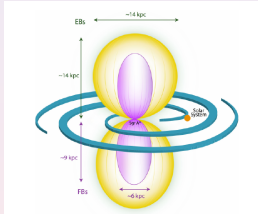


A5: Disentangling cosmic ray signatures in AGN starburst composites



The Astro Projects

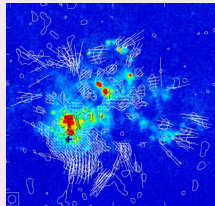
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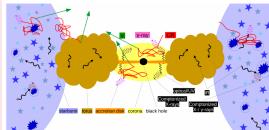
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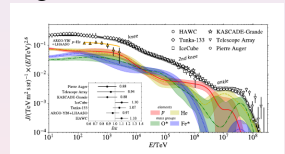
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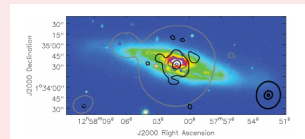
A5: Disentangling cosmic ray signatures in AGN-starburst composites



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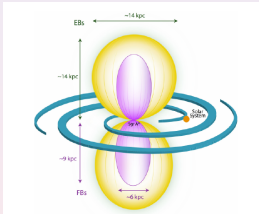


A6: Multimessenger signatures of tidal disruption events

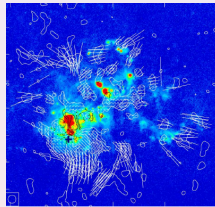


The Astro-Theory Projects

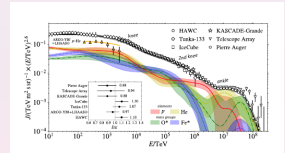
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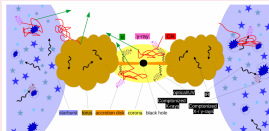
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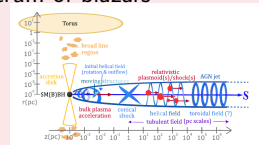
A4: Magnetohydrodynamical halos of starforming galaxies



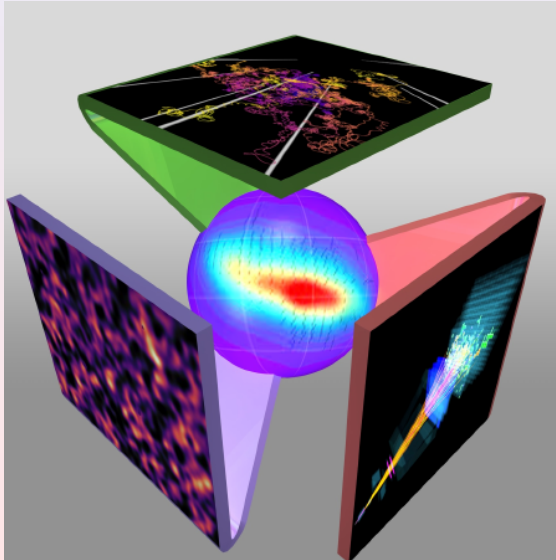
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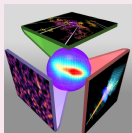
A7: Density-dependence of the temporal structure in the multimessenger spectrum of blazars



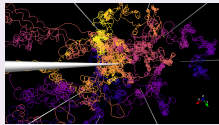
The Astro-Theory Projects and the SFB Logo



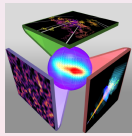
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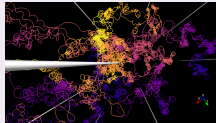
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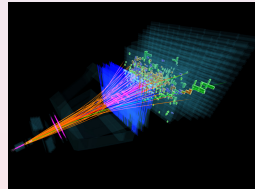
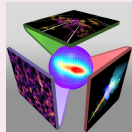
CR transport in
turbulent magne-
tic fields



The Astro-Theory Projects and the SFB Logo

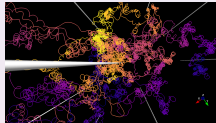


CR transport in
turbulent magne-
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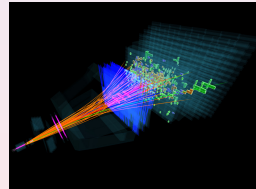
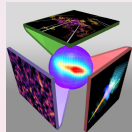


Hadronic interac-
tions in particle
physics and astro-
physics:
CR signatures

The Astro-Theory Projects and the SFB Logo



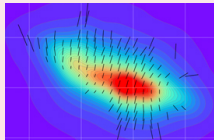
CR transport in turbulent magnetic fields



Hadronic interactions in particle physics and astrophysics:
CR signatures

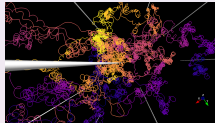


M82: magnetic fields in dynamical halos

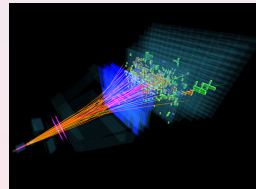
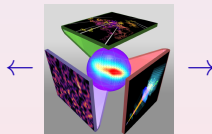
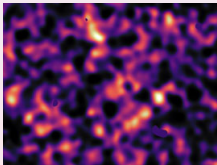


The Astro-Theory Projects and the SFB Logo

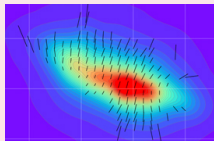
Dark matter in
and around
galaxies



CR transport in
turbulent magne-
tic fields



M82: magnetic
fields in dynam-
ical halos



Hadronic interac-
tions in particle
physics and astro-
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CR signatures

Outline of the Talk

- Introduction (✓)
- CR transport in turbulent magnetic fields
- Turbulent dynamical halos of galaxies
- Dark matter signatures?

Cosmic Ray Transport in Turbulent Magnetic Fields

The CR transport equation

$$\begin{aligned}
 \frac{\partial n}{\partial t} &= \nabla \cdot \left[\overleftrightarrow{D} \nabla n - \vec{U} n \right] && \text{spatial diffusion \& advection} \\
 &+ \frac{\partial}{\partial p} \left[p^2 D_{pp} \frac{\partial}{\partial p} \left(\frac{n}{p^2} \right) \right] && \text{momentum diffusion} \\
 &- \frac{\partial}{\partial p} \left[\left(\frac{dp}{dt} - \frac{p}{3} (\nabla \cdot \vec{U}) \right) n \right] && \text{momentum loss \& adiab. changes} \\
 &- \frac{n}{\tau_f} - \frac{n}{\tau_d} + S && \text{fragmentation \& radioactive decay \& sources}
 \end{aligned}$$

$$\begin{aligned}
 n &= n(\vec{r}, p, t), \quad \overleftrightarrow{D} = \overleftrightarrow{D}(\vec{r}, p, t), \quad \vec{U} = \vec{U}(\vec{r}, t), \quad D_{pp} = D_{pp}(\vec{r}, p, t), \\
 S &= S(\vec{r}, p, t)
 \end{aligned}$$

The CR transport equation

$$\frac{\partial n}{\partial t} = \nabla \cdot \left[\overleftrightarrow{D} \nabla n - \vec{U} n \right] \quad \text{A1, A2, A3, A4, A5, A7}$$

$$+ \frac{\partial}{\partial p} \left[p^2 D_{pp} \frac{\partial}{\partial p} \left(\frac{n}{p^2} \right) \right] \quad \text{A1, A5}$$

$$- \frac{\partial}{\partial p} \left[\left(\frac{dp}{dt} - \frac{p}{3} (\nabla \cdot \vec{U}) \right) n \right] \quad \text{A1, A2, A3, A5, A7}$$

$$- \frac{n}{\tau_f} - \frac{n}{\tau_d} + S \quad \text{A1, A2, A3, A5, A7}$$

$$n = n(\vec{r}, p, t), \quad \overleftrightarrow{D} = \overleftrightarrow{D}(\vec{r}, p, t), \quad \vec{U} = \vec{U}(\vec{r}, t), \quad D_{pp} = D_{pp}(\vec{r}, p, t), \\ S = S(\vec{r}, p, t)$$

Schlickeiser [2002], Becker-Tjus & Merten et al. [2020]

The Central Transport Quantity: The Diffusion Tensor

Study of large-scale transport requires transformation of the diffusion tensor from a (field-aligned) local to a global system:

- Which local system is the most 'natural' one?

The Central Transport Quantity: The Diffusion Tensor

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- The correct choice of the 'perpendicular' directions is important for the case of anisotropic perpendicular diffusion.

The Central Transport Quantity: The Diffusion Tensor

Study of large-scale transport requires transformation of the diffusion tensor from a (field-aligned) local to a global system:

- Which local system is the most 'natural' one?
- One principal direction aligned with magnetic field.
- The correct choice of the 'perpendicular' directions is important for the case of anisotropic perpendicular diffusion.
- The most natural directions are defined with the Frenet-Serret trihedron, i.e. by the curvature and torsion of a magnetic field:

$$\vec{t} = \vec{B}/B \quad ; \quad \vec{n} = (\vec{t} \cdot \nabla) \vec{t}/k \quad ; \quad \vec{b} = \vec{t} \times \vec{n}$$

The Central Transport Quantity: The Diffusion Tensor

The **general transformation** reads: (Effenberger et al. [2012b])

$$\overleftrightarrow{D}_{global} = A^T \overleftrightarrow{D}_{local} A$$

with $A = \begin{pmatrix} n_1 & b_1 & t_1 \\ n_2 & b_2 & t_2 \\ n_3 & b_3 & t_3 \end{pmatrix}$ and $\overleftrightarrow{D}_{local} = \begin{pmatrix} D_{\perp 1} & 0 & 0 \\ 0 & D_{\perp 2} & 0 \\ 0 & 0 & D_{\parallel} \end{pmatrix}$

resulting in

$$D_{11} = D_{\perp 1} n_1^2 + D_{\perp 2} b_1^2 + D_{\parallel} t_1^2$$

$$D_{12} = D_{\perp 1} n_1 n_2 + D_{\perp 2} b_1 b_2 + D_{\parallel} t_1 t_2$$

$$D_{13} = D_{\perp 1} n_1 n_3 + D_{\perp 2} b_1 b_3 + D_{\parallel} t_1 t_3$$

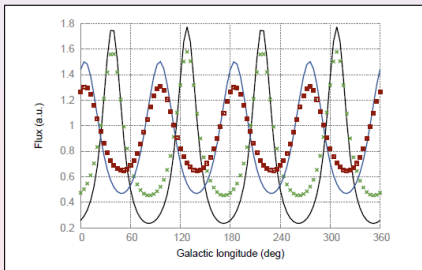
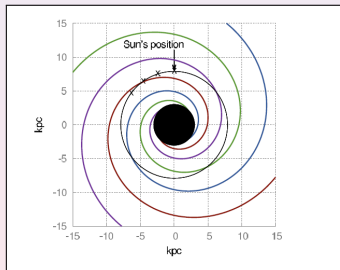
$$D_{22} = D_{\perp 1} n_2^2 + D_{\perp 2} b_2^2 + D_{\parallel} t_2^2$$

$$D_{23} = D_{\perp 1} n_2 n_3 + D_{\perp 2} b_2 b_3 + D_{\parallel} t_2 t_3$$

$$D_{33} = D_{\perp 1} n_3^2 + D_{\perp 2} b_3^2 + D_{\parallel} t_3^2$$

Impact of Anisotropic Diffusion on Galactic Propagation

Orbit of the Sun through the spirally structured galactic cosmic ray (proton) distribution:



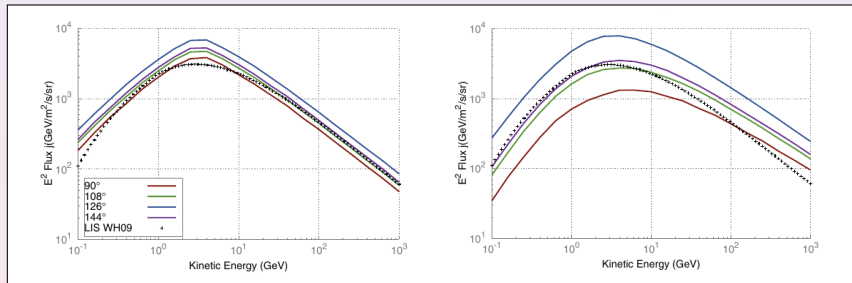
black, green:
1, 100 GeV
at 7.9 kpc

blue, red:
1, 100 GeV
at 5.0 kpc

- stronger intensity variations for anisotropic spatial diffusion
- higher absolute intensities for anisotropic diffusion
(isotropic: $D_{\perp 1} = D_{\perp 2} = D_{\perp} = D_{\parallel}$; anisotropic: $D_{\perp} = 0.01 D_{\parallel}$)

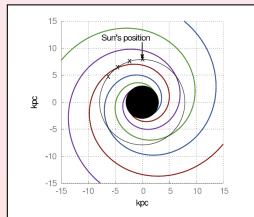
Impact of Anisotropic Diffusion on Galactic Propagation

... and a different local interstellar spectrum:



$$D_{\perp} = D_{\parallel}$$

$$D_{\perp} = 0.01 D_{\parallel}$$



Effenberg et al. [2012b]

Anisotropic Interstellar Diffusion

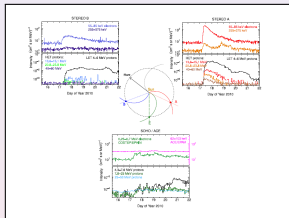
Anisotropic spatial diffusion of cosmic ray protons in the interstellar medium...

- ... causes **stronger intensity variations**;
- ... leads to **higher absolute fluxes** due to a confinement effect;
- ... **influences the local interstellar spectrum**

... at the location of the Sun / Heliosphere in the Galaxy.

The Central Transport Quantity: The Diffusion Tensor

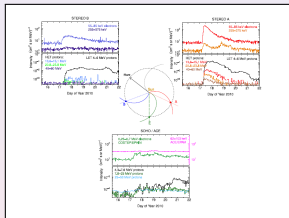
turbulence is crucial for the transport of energetic particles



Dresing et al. [2012]

The Central Transport Quantity: The Diffusion Tensor

turbulence is crucial for the transport of energetic particles



Dresing et al. [2012]

1-comp. model

but:

$$\delta B^2$$

transport process

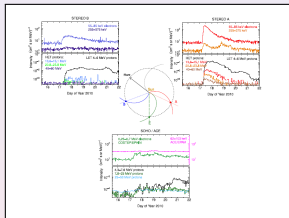
1-comp. model

parallel diffusion D_{\parallel}
 perpendicular diffusion D_{\perp}
 drifts D_A
 momentum diffusion D_{pp}

$$\begin{aligned} &\sim B^2/\delta B^2 \\ &\sim \delta B^2/B^2 \\ &= D_A (\delta B^2/B^2) \\ &= D_{pp} (\delta B^2/B^2) \end{aligned}$$

The Central Transport Quantity: The Diffusion Tensor

turbulence is crucial for the transport of energetic particles



Dresing et al. [2012]

but:

$$\underbrace{\delta B^2}_{1\text{-comp. model}} = \underbrace{\delta B_{2D}^2}_{\substack{\text{'quasi-2D'} \\ \text{low frequency} \\ k \approx k_{\perp}}} + \underbrace{\delta B_{sl}^2}_{\substack{\text{'slab/wave-like'} \\ \text{high frequency} \\ k \approx k_{\parallel}}}$$

transport process

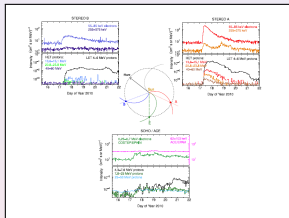
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The Central Transport Quantity: The Diffusion Tensor

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Dresing et al. [2012]

but:

$$1\text{-comp. model} \quad \underbrace{\delta B^2}_{\text{}} = \underbrace{\delta B_{2D}^2}_{\substack{\text{'quasi-2D'} \\ \text{low frequency} \\ k \approx k_{\perp}}} + \underbrace{\delta B_{sl}^2}_{\substack{\text{'slab/wave-like'} \\ \text{high frequency} \\ k \approx k_{\parallel}}}$$

transport process

1-comp. model

2-comp. model

parallel diffusion D_{\parallel}
 perpendicular diffusion D_{\perp}
 drifts D_A
 momentum diffusion D_{pp}

$\sim B^2 / \delta B^2$
 $\sim \delta B^2 / B^2$
 $= D_A (\delta B^2 / B^2)$
 $= D_{pp} (\delta B^2 / B^2)$

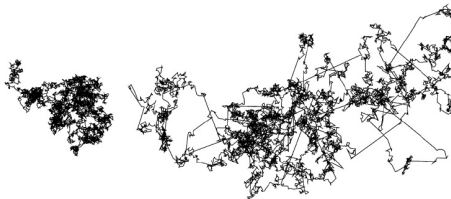
$\sim B^2 / \delta B_{sl}^2$
 $\sim \delta B_{2D}^2 / B^2$
 $= D_A (\delta B_{sl}^2 / B^2, \delta B_{2D}^2 / B^2)$
 $= D_{pp} (\delta B_{sl}^2 / B^2, \delta B_{2D}^2 / B^2)$

What if transport is non-diffusive?

Idea: The particles are not diffusing normally (= 'Gaussian') but anomalously:

$$\langle \Delta x^2 \rangle \sim \begin{cases} t, & \text{Gaussian diffusion} \\ t^\zeta, & \text{anomalous diffusion} \end{cases} \begin{cases} 0 < \zeta < 1, & \text{subdiffusion} \\ 1 < \zeta < 2, & \text{superdiffusion} \end{cases}$$

(sub-)diffusion:
(extended 'waiting times')



super-diffusion:
'Levy flights'

Levy flights \Rightarrow power law distributions

What if transport is non-diffusive?

- ballistic motion: equation of motion

What if transport is non-diffusive?

- ballistic motion: equation of motion

Fokker Planck eq.

- diffusive motion: $\frac{\partial f}{\partial t} = D \frac{\partial^2 f}{\partial x^2}$

What if transport is non-diffusive?

- ballistic motion: equation of motion

(fract.) Fokker Planck eq.

- diffusive motion: $\frac{\partial f}{\partial t} = D \frac{\partial^2 f}{\partial x^2}$

- anom. diff. motion: $\frac{\partial f}{\partial t} = {}_0D_t^{1-\alpha} \left(\hat{D}_\mu \frac{\partial^\mu f}{\partial x^\mu} \right)$

What if transport is non-diffusive?

- ballistic motion: equation of motion

(fract.) Fokker Planck eq. \Leftrightarrow Stoch. Diff. Eqs.

- diffusive motion: $\frac{\partial f}{\partial t} = D \frac{\partial^2 f}{\partial x^2}$ $\Leftrightarrow dX(t) = D^{1/2} dW(t)$

- anom. diff. motion: $\frac{\partial f}{\partial t} = {}_0D_t^{1-\alpha} \left(\hat{D}_\mu \frac{\partial^\mu f}{\partial x^\mu} \right)$ $\Leftrightarrow dX(\tau) = D^{1/\mu} dL_\mu(\tau)$

What if transport is non-diffusive?

- ballistic motion: equation of motion ✓

(fract.) Fokker Planck eq. \Leftrightarrow Stoch. Diff. Eqs. ✓

- diffusive motion: $\frac{\partial f}{\partial t} = D \frac{\partial^2 f}{\partial x^2}$ $\Leftrightarrow dX(t) = D^{1/2} dW(t)$

- anom. diff. motion: $\frac{\partial f}{\partial t} = {}_0D_t^{1-\alpha} \left(\hat{D}_\mu \frac{\partial^\mu f}{\partial x^\mu} \right)$ $\Leftrightarrow dX(\tau) = D^{1/\mu} dL_\mu(\tau)$

Simulation framework **CRPropa** (to be used extensively in the SFB) offers a 'unifying' treatment regarding transport and energy regimes

→ see talks by P. Reichherzer and J. Dörner

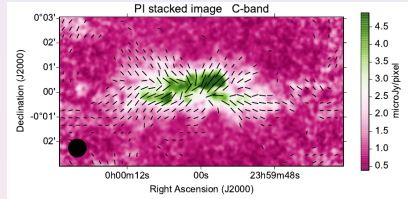
Turbulent Dynamical Halos

Dynamical Halos



polarized intensity
of NGC 4217

→ see talk by M. Stein



stacked images of 28
CHANG-ES galaxies

MHD Description

- large-scale equations:

$$\partial_t \rho + \nabla \cdot (\rho \vec{U}) = 0$$

$$\partial_t (\rho \vec{U}) + \nabla \cdot \left[\rho \vec{U} \vec{U} + \left(\rho + \frac{|\vec{B}|^2}{2} + p_w \right) \mathbb{1} - \left(1 + \frac{\sigma_D \rho Z^2}{2B^2} \right) \vec{B} \vec{B} \right] = -\rho \vec{g}$$

$$\begin{aligned} \partial_t \mathbf{e} + \nabla \cdot \left[\mathbf{e} \vec{U} + \left(\rho + \frac{|\vec{B}|^2}{2} \right) \vec{U} - (\vec{U} \cdot \vec{B}) \vec{B} \right] + \vec{q}_H - \frac{\rho H_c}{2} \vec{V}_A \\ = -\rho \vec{U} \cdot \vec{g} - \vec{U} \cdot \nabla p_w - (\vec{V}_A \cdot \nabla \rho) \frac{H_c}{2} + \frac{\rho Z^3 f}{2\lambda} \\ + \vec{U} \cdot (\vec{B} \cdot \nabla) \left[\frac{\sigma_D \rho Z^2}{2B^2} \vec{B} \right] - \rho \vec{V}_A \cdot \nabla H_c \end{aligned}$$

$$\partial_t \vec{B} + \nabla \cdot (\vec{U} \vec{B} - \vec{B} \vec{U}) = \vec{0}$$

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$$\partial_t H_c + \nabla \cdot (\vec{U} H_c + \vec{V}_A Z^2) = \frac{H_c}{2} \nabla \cdot \vec{U} + 2\vec{V}_A \cdot \nabla Z^2 + Z^2 \sigma_D \nabla \cdot \vec{V}_A - \frac{\alpha Z^3 f^-}{\lambda}$$

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$$\rho_w = \delta B^2 = (1 + \sigma_D) \rho Z^2 / 4 \Rightarrow \overset{\leftrightarrow}{D}$$

$$\partial_t Z^2 + \nabla \cdot (\vec{U} Z^2 + \vec{V}_A H_c) = \frac{Z^2(1 - \sigma_D)}{2} \nabla \cdot \vec{U} + 2\vec{V}_A \cdot \nabla H_c + \frac{Z^2 \sigma_D}{|\vec{B}|^2} \vec{B} \cdot (\vec{B} \cdot \nabla) \vec{U} - \frac{\alpha Z^3 f^+}{\lambda} + S$$

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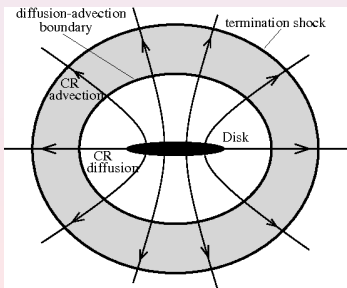
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MHD Simulation Output

- large-scale eqs: velocity field, magnetic field, ... → A1, A2
- small-scale eqs: turbulence... → F1 → (modified) diffusion tensor
→ see talk by F. Effenberger

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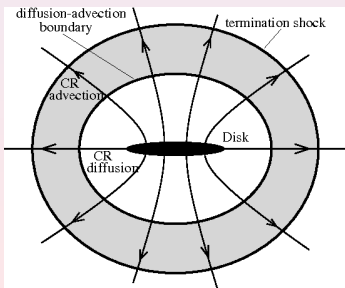


→ see talk by A. Käpä

*Jokipii & Morfill [1985],
Völk & Ziraqashvili [2003],
Merten et al. [2018]*

MHD Simulation Output

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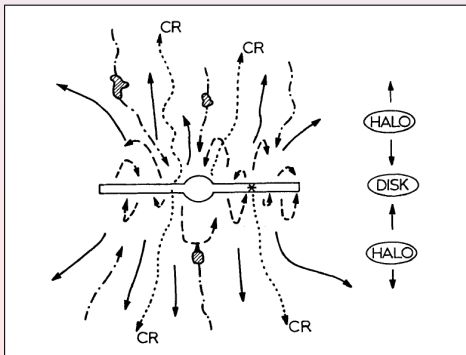
*Jokipii & Morfill [1985],
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Merten et al. [2018]*

In the MHD code **Cronos** the large- and small-scale equations are implemented already (originally for heliospheric applications)

Dark Matter

Consequences of Dark Matter

- A1: Dark matter signature from Galactic center?
- A2: Dark matter signatures from dwarf galaxies?
- A4: Impact on halo dynamics?
 - global winds, partial winds, fountains, ... ?

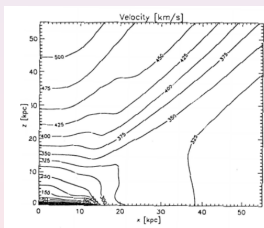


Völk et al. [1989]

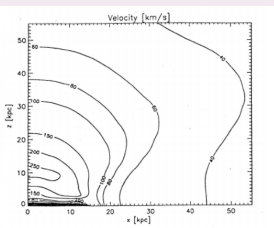
Consequences of Dark Matter

Adding a (spherical/axisymmetric/triaxial) dark halo potential (determined in F6) may also result in *breeze solutions*:

Wind



Breeze

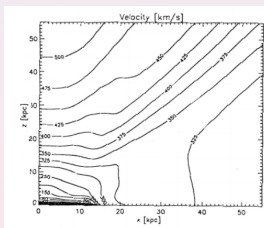


Fichtner & Vormbrock [1997], Taylor & Giacinti [2017]

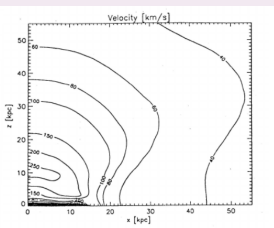
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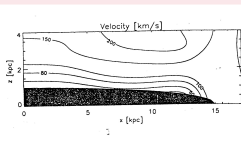
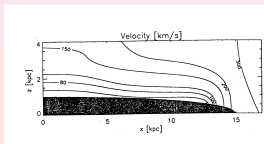
Wind



Breeze



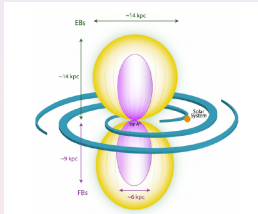
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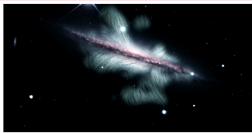
Résumé

Well... instead of a Résumé: Check the next Annual Meeting ☺

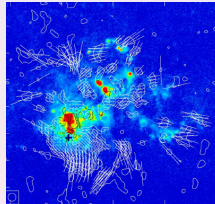
A1: Multimessenger signatures of Gal. CR transport



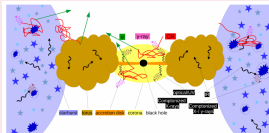
A4: Magnetohydrodynamical halos of starforming galaxies



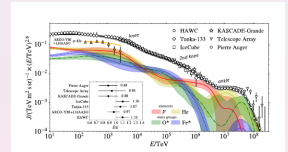
A2: CR signatures in dwarf galaxies



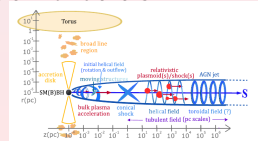
A5: Disentangling cosmic ray signatures in AGN-starburst composites



A3: CR transport in the transition region from Galactic to extragalactic origin



A7: Density-dependence of the temporal structure in the multimessenger spectrum of blazars



Well... instead of a Résumé: Check the next Annual Meeting ☺

A1: Multimessenger signatures of Gal. CR transport

A2: CR signatures in dwarf galaxies

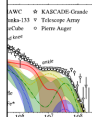
A3: CR transport in the transition region from Galactic to extragalactic

Modelling of

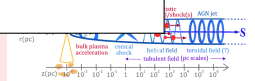
- thermal plasma dynamics
- cosmic ray transport
- turbulence
- multimessenger signatures

not new in all aspects but new in this coordinated fashion

A4: M...
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galaxie



dependence of
structure in
galaxy spec-

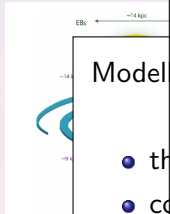


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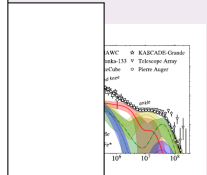
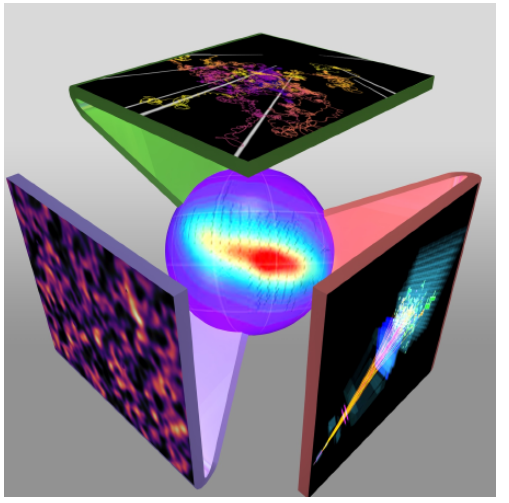
A3: CR transport in the transition region from Galactic to extragalactic



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A4: M...
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